Supplemental Materials for "Inverse design of phononic crystal with desired transmission via a gradient-descent approach"

Yuhang Wei(魏宇航) and Dahai He(贺达海)*

Department of Physics and Jiujiang Research Institute, Xiamen University, Xiamen 361005, China

THE TRANSMISSION IN TERMS OF NORMAL MODES

From the motion equation $M\ddot{X} + KX = 0$, we know that the amplitude X of the normal mode satisfies $\omega^2 MX = KX$. Let $U = M^{\frac{1}{2}}X$, then we can get $\omega^2 U = M^{-\frac{1}{2}}KM^{-\frac{1}{2}}U$ which satisfies

$$F\Omega^2 F^{\top} = M^{-\frac{1}{2}} K M^{-\frac{1}{2}}, \tag{S1}$$

where eigenmatrix $F_{ji} = e_j^i$ and eigenfrequencies $\Omega_{ii} = \omega_i$. Define $R = \Omega^2 - \omega^2 I$ and $Z' = Z + \Sigma(\omega) = K - \omega^2 M$ where $Z'_{1N}^* = Z_{1N}^*$ and I is identity matrix, then we obtain $Z' = M^{\frac{1}{2}}FRF^{\top}M^{\frac{1}{2}}$. Thus $Z_{1N}^* = \text{Det}(MR)(M^{-\frac{1}{2}}FR^{-1}F^{\top}M^{-\frac{1}{2}})_{1N}$. Then one can find

$$Z_{1N}^* = \frac{\text{Det}(MR)}{\sqrt{\gamma_{\rm L}\gamma_{\rm R}}} S_1 = 0, \tag{S2}$$

where S_1 is given by

$$S_1 = \sum_{i=1}^{N} \frac{E_1^i E_N^i}{\omega_i^2 - \omega^2}.$$
 (S3)

Next, we simplify Det(Z). Due to the property of determinant, one has $\text{Det}(Z) = \text{Det}(MF^{\top}M^{-\frac{1}{2}}ZM^{-\frac{1}{2}}F)$ and

$$F^{\top}M^{-\frac{1}{2}}ZM^{-\frac{1}{2}}F = R + A + B,$$
(S4)

where $A = -i \frac{\gamma \mathbf{L}}{m_1} F_{(1,:)}^{\top} F_{(1,:)}, B = -i \frac{\gamma \mathbf{R}}{m_N} F_{(N,:)}^{\top} F_{(N,:)}$. The $F_{(i,:)}$ means taking the *i*th row of matrix F as a vector. Since R is a diagonal matrix and the rank of matrices A and B is 1, one can find

$$Det(R + A + B) = Det(R) \left[1 + \sum_{i=1}^{N} \frac{A_{ii} + B_{ii}}{R_{ii}} + \sum_{1=i (S5)$$

It results in

$$Det(Z) = Det(MR)(1 - i\omega S_3 - \omega^2 S_2),$$
(S6)

where S_2 and S_3 is given by

$$S_{2} = \sum_{1=i

$$S_{3} = \sum_{i=1}^{N} \frac{(E_{1}^{i})^{2} + (E_{N}^{i})^{2}}{\omega_{i}^{2} - \omega^{2}}.$$
(S7)$$

So far we have obtained Z_{1N}^* and Det(Z). Finally, in terms of $G_{1N} = Z_{1N}^*/\text{Det}(Z)$ and $\mathcal{T}(\omega) = 4\gamma_L \gamma_R \omega^2 |G_{1N}|^2$, one can get the transmission with respect to the normal mode as given by

$$\mathcal{T}(\omega) = \frac{4\omega^2 S_1^2}{(1 - \omega^2 S_2)^2 + \omega^2 S_3^2}.$$
(S8)

* dhe@xmu.edu.cn