

# Supporting Information for Cobalt-Dimer Nitrides: A Potential Novel Family of High Temperature Superconductors

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## COMPUTATIONAL METHODS

We employ the Vienna *ab initio* simulation package (VASP) code[1] with the projector augmented wave (PAW) method[2] to perform DFT calculations. The Perdew-Burke-Ernzerhof (PBE)[3] exchange-correlation functional is used in our calculations. The kinetic energy cutoff is set to be 600 eV for expanding the wave functions into a plane-wave basis in VASP calculations. We use the primitive cell of BaCo<sub>2</sub>N<sub>2</sub> to calculate its electronic structure. The crystal structure is fully relaxed while the force convergence criterion is 0.005 eV/Å. The energy convergence criterion is 10<sup>-7</sup> eV and the  $\Gamma$ -centered  $\mathbf{k}$ -mesh is  $12 \times 12 \times 12$ . The Co<sub>2</sub>N<sub>2</sub> monolayer is extracted from the relaxed bulk BaCo<sub>2</sub>N<sub>2</sub> with 40Å thick vacuum layer and its  $\Gamma$ -centered  $\mathbf{k}$ -mesh is  $16 \times 16 \times 1$ .

We employ Wannier90[4, 5] to calculate maximally localized Wannier functions in BaCo<sub>2</sub>N<sub>2</sub>, which perfectly reproduces DFT-calculated band structure around the Fermi level. The initial projectors are Co's *d*-orbitals with local symmetric  $X - Y$  coordinate.

In the study of the magnetism of BaCo<sub>2</sub>N<sub>2</sub>, the DFT+ $U$  method in the formulation of Liechtenstein et al. [8] is employed to describe the associated electron-electron correlation effect and  $J$  is fixed as  $\frac{U}{4}$ . To estimate the magnetic exchange between two NN dimers, the energies of different magnetic states are calculated with the relaxed structure of paramagnetic BaCo<sub>2</sub>N<sub>2</sub>.

## EXPLICIT FORM OF 4-BAND TB MODEL

The 4-band effective tight-binding (TB) model before unfolding with the basis of the  $d_{XY}$  and  $d_{YZ}$  orbitals is given by

$$H_t^{4-band} = \sum_{k\alpha\beta ij} \varepsilon_k^{\alpha\beta ij} d_{k\alpha i}^\dagger d_{k\beta j} + \sum_{ki} e_\alpha d_{k\alpha i}^\dagger d_{k\alpha i}, \quad (1)$$

where  $\alpha/\beta$  is the orbital index,  $i/j$  is the site index and  $e_\alpha$  denotes the crystal field energy of each orbital. The corresponding expressions for the hopping energies in Eq. (1) are given by

$$\varepsilon_k^{\alpha\alpha AA} = \varepsilon_k^{\alpha\alpha BB} = 2t'_\alpha (\cos k_x + \cos k_y); \quad (2)$$

$$\varepsilon_k^{\alpha\alpha AB} = \varepsilon_k^{\alpha\alpha BA} = 4t_\alpha \cos \frac{k_x}{2} \cos \frac{k_y}{2}. \quad (3)$$

As mentioned in our main text, we can transfer the  $4 \times 4$  TB model into a block-diagonalized matrix with using the glide symmetry:

$$H^{eff}(\mathbf{k}) = U H_t^{4-band} U^\dagger = \begin{pmatrix} H_t(\mathbf{k}) & 0 \\ 0 & H_t(\mathbf{k}+\mathbf{Q}) \end{pmatrix}, \quad (4)$$

here  $H_k$  is the effective two-band model in our main text and  $\mathbf{Q} = (\pi, \pi)$  is the folding vector. The unitary transformation  $U$  is defined as

$$U = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 0 & 1 & 0 \\ 0 & 1 & 0 & 1 \\ 1 & 0 & -1 & 0 \\ 0 & 1 & 0 & -1 \end{pmatrix}, \quad (5)$$

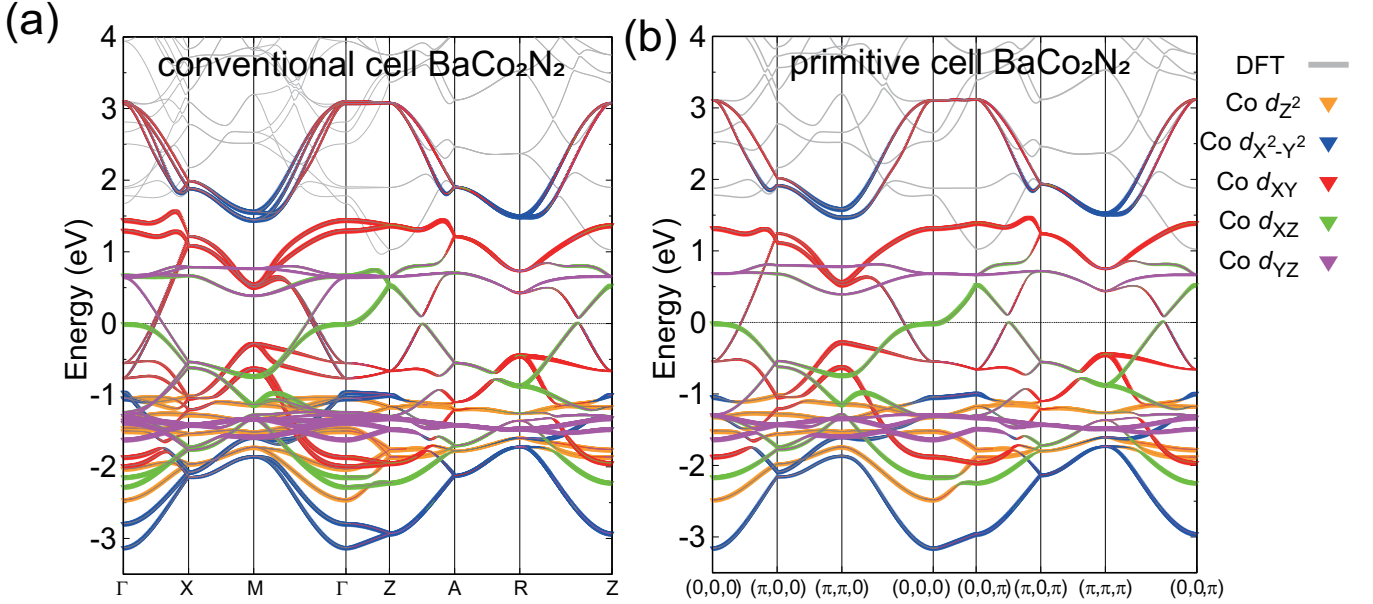


FIG. S1: DFT-calculated band structures of  $\text{BaCo}_2\text{N}_2$  and projected weight of  $d$ -orbital-like WFs in its symmetric local coordinate ( $X$ - $Y$  coordinate) with (a) the conventional cell and (b) the primitive cell. The coordinates in the abscissa of (b) are corresponding to the high symmetric  $\mathbf{k}$ -points of the tetragonal conventional cell's Brillouin zone.

which is similar to that in  $\text{FeSe}$ [6] and  $\text{BaCuS}_2$ [7].

#### BAND STRUCTURE OF $\text{BaCo}_2\text{N}_2$

Fig.S1 shows the orbital-resolved band structure from DFT calculations of  $\text{BaCo}_2\text{N}_2$ . The N  $p$  orbitals, absent in the figure, are located around 3.5 eV below the Fermi level, while Co  $3d$  orbitals dominate from -3 eV to 3 eV. Despite moderate inter-orbital coupling, low-energy bonding states and high-energy anti-bonding states for each Co  $3d$  orbital can be identified. We also notice that the dispersion of band structure around the Fermi level is similar with that in  $\text{Co}_2\text{N}_2$  monolayer, which is consistent with our previous qualitative analysis from symmetry.

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- [1] G. Kresse and J. Furthmüller, Phys. Rev. B 54, 11169 (1996).
- [2] G. Kresse and D. Joubert, Phys. Rev. B 59, 1758 (1999).
- [3] J. P. Perdew, K. Burke and M. Ernzerhof, Phys. Rev. Lett. 77, 3865 (1996).
- [4] A. A. Mostofi, J. R. Yates, Y.-S. Lee, I. Souza, D. Vanderbilt, and N. Marzari, Comput. Phys. Commun. 178, 685 (2008).
- [5] N. Marzari, A. A. Mostofi, J. R. Yates, I. Souza, and D. Vanderbilt, Rev. Mod. Phys. 84,1419 (2012).
- [6] K. Jiang, J. Hu, H. Ding, and Z. Wang, Phys. Rev. B **93**, 115138 (2016).
- [7] Y. Gu, X. Wu, K. Jiang, and J. Hu, Chin. Phys. Lett. **38**, 017501 (2021).
- [8] A. Liechtenstein, V. I. Anisimov, and J. Zaanen, Phys. Rev. B 52, R5467 (1995).