

# Supporting Information: Directional Design of Two-Dimensional SnSe Materials with High Thermopower Based on the Pareto Optimization

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## I. CORRELATION BETWEEN BAND EFFECTIVE MASS AND DENSITY OF STATES EFFECTIVE MASS, AND SEEBECK COEFFICIENT FOR 3D AND 2D MATERIALS

Here, we show the different correlations between Seebeck coefficient and density of states (DOS) effective mass ( $m_{DOS}^*$ ) for 3D and 2D materials. Meanwhile, we can also obtain the correlations between band effective mass ( $m_{band}^*$ ) and density of states (DOS) effective mass ( $m_{DOS}^*$ ) for 3D and 2D materials.

According to the Onsager reciprocal relations of the charge and heat currents [1]:  $j_Q = L^{(0)} \cdot \mathbf{E} + \frac{L^{(1)}}{qT} \cdot (-\nabla T)$ ,  $j_Q = \frac{L^{(1)}}{q} \cdot \mathbf{E} + \frac{L^{(2)}}{q^2 T} \cdot (-\nabla T)$ , where  $L^{(\alpha)} = q^2 \int \sigma(E)(E - E_F)^\alpha (-\frac{\partial f(E)}{\partial E}) dE$ , we can obtain the electronic conductivity and the Seebeck coefficient under the parabolic band model approximation:  $\sigma_0 = L^{(0)}$ ,  $S = \frac{1}{qT} \frac{L^{(1)}}{L^{(0)}}$ , where the  $q$ ,  $\mathbf{E}$ ,  $E$ ,  $T$ ,  $E_F$  and  $f(E)$  are the charge, electric field, energy, temperature, Fermi energy and Fermi-Dirac distribution function, respectively.

When only considering the electronic transport near the Fermi level, after performing a Sommerfeld expansion, we can obtain the Mott's relations [2]:  $S = \frac{\pi^2 k_B^2 T}{3q} \left. \frac{d[\ln(\sigma(E))]}{dE} \right|_{E=E_F}$ . We can find that the DOS will make an important effect on the Seebeck coefficient.

The  $\sigma(E)$  is the energy-dependent electronic conductivity, expressed as:  $\sigma(E) = \tau(E)v^2(E)G(E)$ , where the  $\tau(E)$ ,  $v(E)$  and  $G(E)$  are the electronic relaxation time, group velocity of carriers and DOS of multiple degenerate bands. The DOS of multiple degenerate bands,  $G(E)$  is determined by the DOS of single band:  $G(E) = N_v g(E)$ , where the  $N_v$  and  $g(E)$  are band degeneracy and the DOS of single band respectively. If only considering the acoustic phonon scattering, we can get:  $\tau(E) \propto E^{-1/2}$  [3]. According to the free electron approximation, the correlation between group velocity of carriers and energy is  $v(E) \propto E^{1/2}$ . Here we will give the DOS expression for 3D and 2D systems.

For 3D system, the DOS of single band model is:  $g_{3D}(E) = \frac{1}{2\pi^2} (\frac{2m_{band}^*}{\hbar^2})^{3/2} E^{1/2}$  where  $\hbar$  is reduced Planck constant. If considering the multiple degenerate bands model, the DOS of multiple degenerate bands model,  $G_{3D}(E)$  is:  $G_{3D}(E) = \frac{1}{2\pi^2} (\frac{2m_{DOS}^*}{\hbar^2})^{3/2} E^{1/2} = N_v g_{3D}(E)$ . Then we can obtain that:  $G_{3D}(E) = \frac{2}{2\pi^2} (\frac{2N_v^{2/3} m_{band}^*}{\hbar^2})^{3/2} E^{1/2}$ . Therefore, we can find the correlation among the DOS effective mass, band degeneracy and single band effective mass for the 3D system [4, 5]:  $m_{DOS}^* = N_v^{2/3} m_{band}^*$ .

Similarly, for 2D system, the DOS of single band model is:  $g_{2D}(E) = \frac{2}{(2\pi)^2} \frac{2\pi}{\hbar} m_{band}^*$ . If considering the multiple degenerate bands model, the DOS of multiple degenerate bands model,  $G_{2D}(E)$  is:  $G_{2D}(E) = \frac{2}{(2\pi)^2} \frac{2\pi}{\hbar} m_{DOS}^* = N_v g_{2D}(E)$ . Hence, we also can obtain that:  $G_{2D}(E) = \frac{2}{(2\pi)^2} \frac{2\pi}{\hbar} N_v m_{band}^*$ . We can also find the correlation among the DOS effective mass, band degeneracy and single band effective mass for the 2D system:  $m_{DOS}^* = N_v m_{band}^*$ . Then, we give the formulas of the Seebeck coefficient for 3D and 2D systems.

For 3D system, the energy-dependent electronic conductivity:  $\sigma_{3D}(E) \propto E$ , we could obtain the Seebeck coefficient:

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$$\begin{aligned}
S_{3D} &= \frac{\pi^2 k_B^2 T}{3q} \left. \frac{d[\ln(\sigma(E))]}{dE} \right|_{E=E_F} \\
&= \frac{\pi^2 k_B^2 T}{3q} \frac{1}{E_F} \\
&= \frac{\pi^2 k_B^2 T}{3q} \frac{2m_{DOS,F}^*}{\hbar^2 k_{3D,F}^2} \\
&= \frac{2\pi^2 k_B^2 T}{3q \hbar^2} \left( \frac{1}{3\pi^2 n} \right)^{2/3} m_{DOS,F}^*
\end{aligned} \tag{1}$$

where  $m_{DOS,F}^* = \frac{\hbar^2 k_F^2}{2E_F}$  is the DOS effective mass at the Fermi energy level,  $k_F$  is the Fermi radius, the  $k_{3D,F} = (3\pi^2 n)^{1/3}$  is the Fermi radius for 3D system, and  $n$  is carrier concentration.

Similarly, for 2D system, the energy-dependent electronic conductivity:  $\sigma_{2D}(E) \propto E^{1/2}$ , we could obtain the Seebeck coefficient:

$$\begin{aligned}
S_{2D} &= \frac{\pi^2 k_B^2 T}{3q} \left. \frac{d[\ln(\sigma(E))]}{dE} \right|_{E=E_F} \\
&= \frac{\pi^2 k_B^2 T}{3q} \frac{1}{2} \frac{1}{E_F} \\
&= \frac{\pi^2 k_B^2 T}{3q} \frac{1}{2} \frac{2m_{DOS,F}^*}{\hbar^2 k_{2D,F}^2} \\
&= \frac{\pi k_B^2 T}{6qn \hbar^2} m_{DOS,F}^*
\end{aligned} \tag{2}$$

where the  $k_{2D,F} = (2\pi n)^{1/2}$  is the Fermi radius for 2D system.

## II. THE THERMOPOWER ALONG X AND Y DIRECTIONS

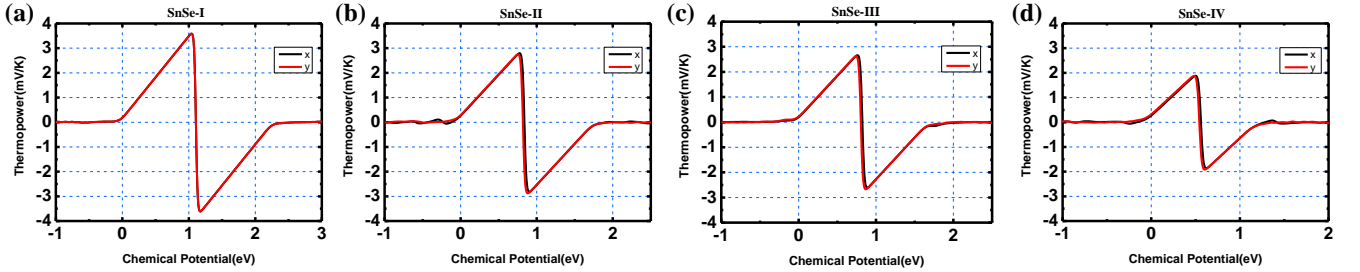


FIG. 1: There is thermopower of the four structures along x(black line) and y(red line) directions at 300K.

### III. THE BILAYER STRUCTURES ON THE FIRST PARETO FRONT

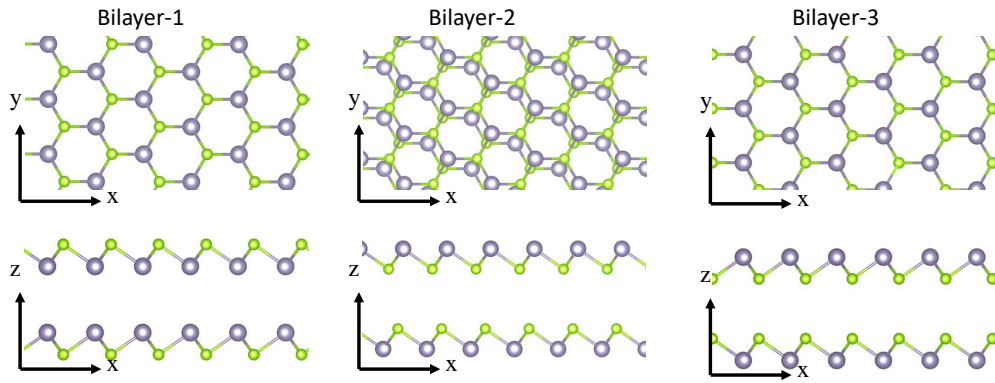


FIG. 2: There are three bilayer structures of 2D SnSe with the top and side views, which are blue hollow hexagonal symbols, as shown on the first Pareto Front.

TABLE I: Properties of three bilayer structures of 2D SnSe. There are space group, the values of lattice parameters  $a$  and  $b$ , free energy per atom and the values of thermopower, respectively.

Materials	Space group	$a(\text{\AA})$	$b(\text{\AA})$	$E(\text{eV/atom})$	$S(\text{mV/K})$
Bilayer-1	P-6m2	3.920	3.920	-4.081	2.829
Bilayer-2	P-1	3.913	3.915	-4.080	2.965
Bilayer-3	P-6m2	3.910	3.910	-4.079	2.968

## IV. STRUCTURES ON THE SECOND PARETO FRONT AND THIRD PARETO FRONT

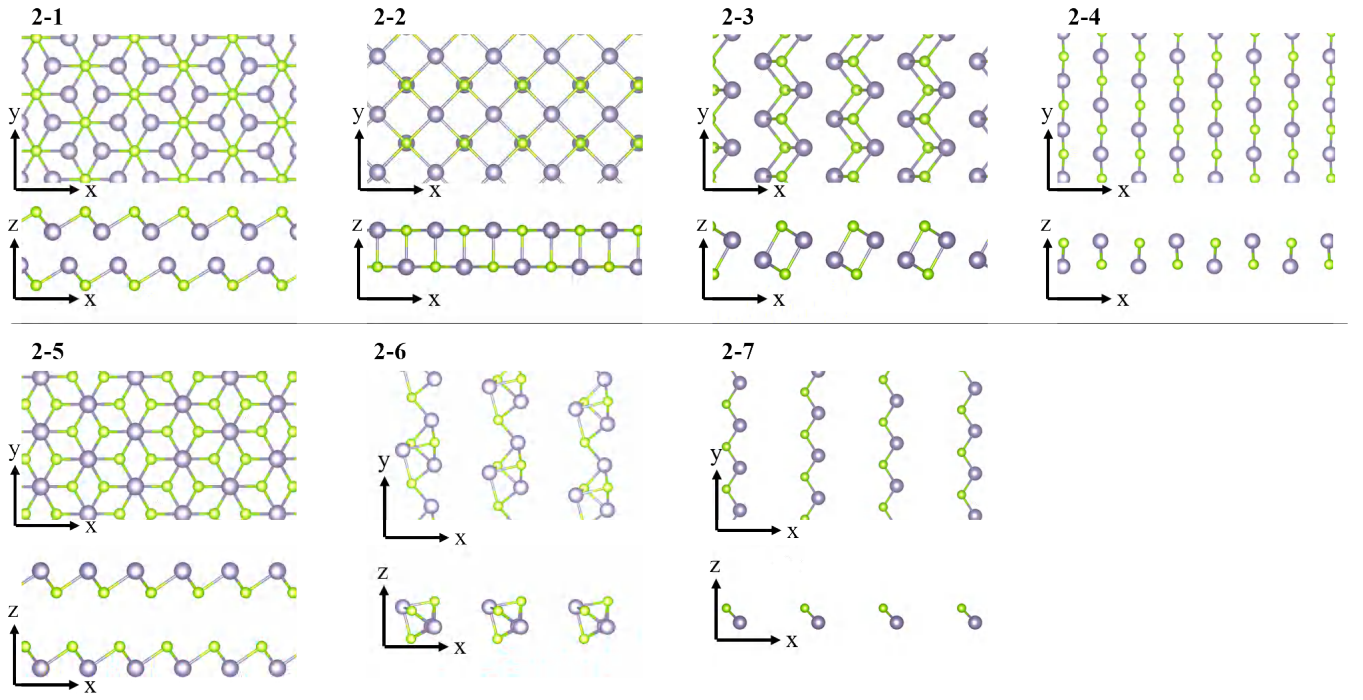


FIG. 3: Structures of 2D SnSe on the second Pareto Front are sorted by the free energy.

TABLE II: Properties of 2D SnSe structures on the second Pareto Front. There are space group, the values of lattice parameters  $a$  and  $b$ , free energy per atom and the values of thermopower, respectively.

Materials	Space group	$a(\text{\AA})$	$b(\text{\AA})$	$E(\text{eV/atom})$	$S(\text{mV/K})$
2-1	P-3m1	4.094	4.094	-4.137	0.887
2-2	Cmcm	4.320	4.320	-4.135	1.320
2-3	P-1	4.112	5.109	-4.105	1.555
2-4	Ama2	3.937	6.148	-4.103	2.719
2-5	C2/m	6.795	3.918	-4.080	2.906
2-6	P1	7.031	7.329	-3.990	2.914
2-7	P1	3.800	7.134	-3.794	2.971

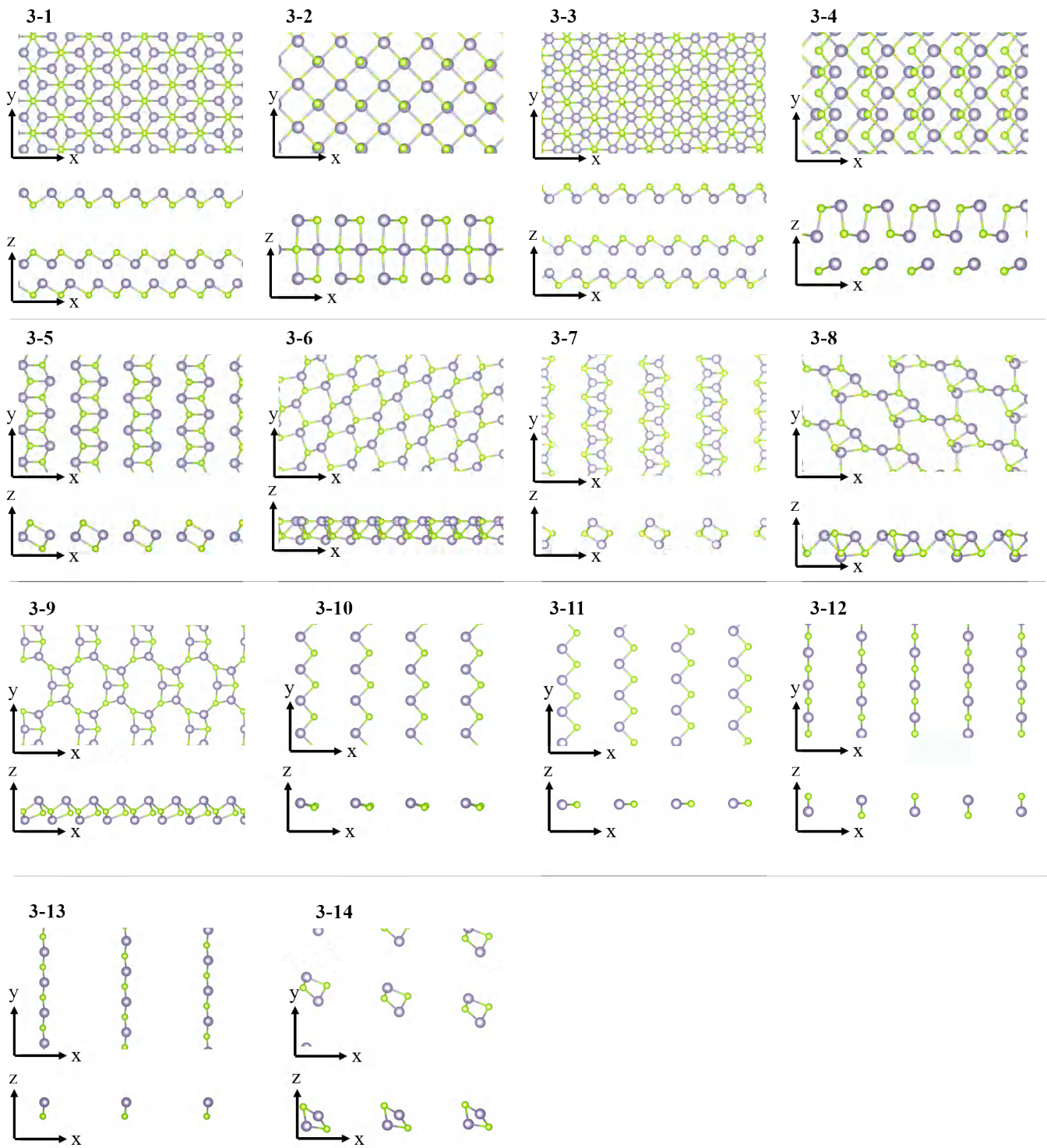


FIG. 4: Structures of 2D SnSe on the third Pareto Front are sorted by the free energy.

TABLE III: Properties of 2D SnSe structures on the third Pareto Front. There are space group, the values of lattice parameters  $a$  and  $b$ , free energy per atom and the values of thermopower, respectively.

Materials	Space group	$a(\text{\AA})$	$b(\text{\AA})$	$E(\text{eV/atom})$	$S(\text{mV/K})$
3-1	P3m1	4.039	4.039	-4.110	0.961
3-2	Cm	6.154	6.039	-4.107	1.003
3-3	Cm	6.825	3.989	-4.096	1.054
3-4	Pm	4.532	4.134	-4.077	1.359
3-5	P-1	3.911	6.410	-4.071	2.178
3-6	P1	6.516	6.501	-4.047	2.279
3-7	P-1	3.955	7.964	-4.034	2.720
3-8	P1	6.109	7.060	-4.027	2.762
3-9	P1	7.224	7.232	-3.986	2.775
3-10	Pm	6.594	7.615	-3.801	2.778
3-11	Amm2	3.793	13.768	-3.800	2.813
3-12	Pmmn	3.789	12.448	-3.797	2.894
3-13	P1	3.816	10.770	-3.787	2.963
3-14	P1	8.306	9.612	-3.738	2.970

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